

Vibration-rotation kinetic energy operators and geometric algebra

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Outline

- 1 Introduction
- 2 KEOs by GA
- 3 Conclusions

Molecular energy levels

- In 1920's it was observed that molecular energy levels are quantized
- IR spectrum shows band structure

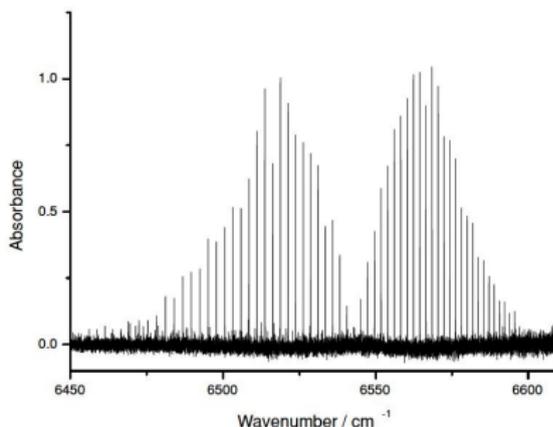


Figure: A part of absorption spectrum of $^{12}\text{C}^{13}\text{CH}$, recorded in the Laboratory of physical chemistry, University of Helsinki.

Molecular energy levels (cntnd.)

- To explain observed spectras (or to predict energy level structure), we need to solve the Schrödinger equation

$$H\Psi = E\Psi \quad (1)$$

- **Eigenfunction** Ψ and the **eigenvalue** E are unknown.
- **Hamiltonian** $H = T + V$ is known

Vibration-rotation SY

- Born-Oppenheimer approximation assumed!
 - Molecule in an electronic eigenstate.
 - Potential V obtained by fitting the electronic energies for given values of nuclear coordinates.
- Only nuclear coordinates included.

Vibration-rotation KEO

- Vibration-rotation Kinetic energy operator (KEO) given by

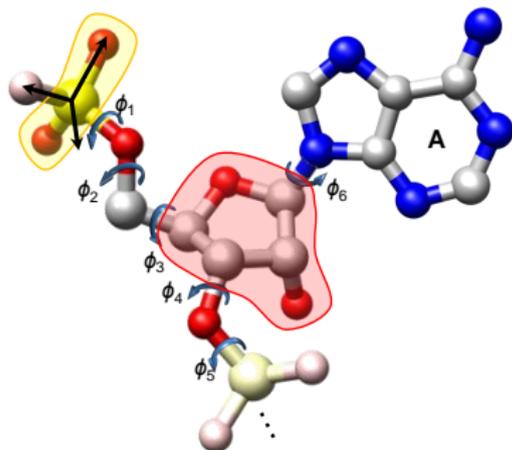
$$T = -\frac{\hbar^2}{2} \sum_{\alpha}^N \nabla_{\mathbf{x}_{\alpha}}^2 \quad (2)$$

- The **Laplacian** $\nabla_{\mathbf{x}_{\alpha}}^2 = \nabla_{\mathbf{x}_{\alpha}} \cdot \nabla_{\mathbf{x}_{\alpha}}$ is the square of the **gradient operator** $\nabla_{\mathbf{x}_{\alpha}}$, which operates on the position vector \mathbf{x}_{α} of the atom α .

Vibration-rotation KEO (cntd.)

- The degrees of freedom that a polyatomic molecule possess are classified into three different types:
 - **Position** (of the center-of mass) – 3 cartesian coordinates.
 - Translation.
 - **Orientation** (of the body-frame w.r.t. constant lab-frame) – 3 Euler angles.
 - Rotation.
 - **Shape** – N_s ($\leq 3N - 6$) internal coordinates.
 - Deformation, vibration.

Vibration-rotation KEO (cntd.)



Vibration-rotation KEO (cntd.)

- In internal coordinates, the **vibration-rotation** part of KEO reads as

$$\begin{aligned} T = & \frac{1}{2} \sum_{ij}^{3N-6} p_i^\dagger g^{(s_i s_j)} p_j + \frac{1}{2} \sum_i^{3N-6} \sum_j^3 \left(p_i^\dagger g^{(s_i l_j)} l_j + l_j g^{(s_i l_j)} p_i \right) \\ & + \frac{1}{2} \sum_{ij}^3 l_i^\dagger g^{(l_i l_j)} l_j \end{aligned} \quad (3)$$

Vibration-rotation KEO (cntd.)

- $p_i = -i\hbar\partial/\partial s_i$ and $p_i^\dagger = -i\hbar[\partial/\partial s_i + (\partial J/\partial s_i)/J]$
 - J is the **Jacobian**.
- Angular momentum operator is self-adjoint, i.e., $l_i^\dagger = l_i$.
- $g^{(s_i s_j)}$ is the **vibrational**, $g^{(s_i l_j)}$ is the **Coriolis**, and $g^{(l_i l_j)}$ is the **rotational** element of the **contravariant metric tensor**.
 - Functions of the shape coordinates only.
 - Euler angles do **not** appear in metric tensor elements!

Vibration-rotation KEO (cntd.)

- Q: How to obtain KEO and related quantities?
- A: By **Geometric algebra** (= Real Clifford algebra with geometric interpretation).

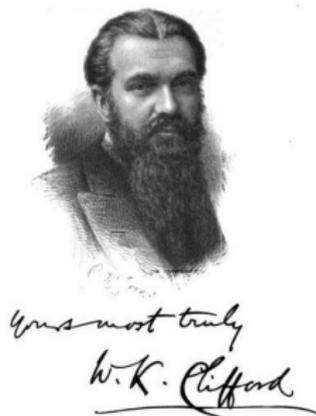


Figure: William Kingdon Clifford (1845-1879)

Geometric (Clifford) algebra

- Algebra of **Euclidean displacements** (vectors)
 - $\mathbf{a}(\mathbf{bc}) = (\mathbf{ab})\mathbf{c} = \mathbf{abc}$
 - $\mathbf{a}(\mathbf{b} + \mathbf{c}) = \mathbf{ab} + \mathbf{ac}$
 - $(\mathbf{b} + \mathbf{c})\mathbf{a} = \mathbf{ba} + \mathbf{ca}$
- Reflects directly the geometric properties:
 - $\mathbf{a} \parallel \mathbf{b} \Leftrightarrow \mathbf{ab} \in \mathbb{R}$ (especially, $\mathbf{a}^2 \equiv |\mathbf{a}|^2$).
 - $\mathbf{a} \perp \mathbf{b} \Leftrightarrow \mathbf{ab} = -\mathbf{ba}$

Geometric (Clifford) algebra (cntnd.)

- Vectors used in calculations, not their components!

Example

For example,

$$\mathbf{a} \rightarrow \mathbf{a}' = -\mathbf{u}\mathbf{a}\mathbf{u}$$

is the **reflection** of \mathbf{a} along the (unit) vector \mathbf{u} (it reverses the part of \mathbf{a} , which is perpendicular to \mathbf{u} and does nothing to the parallel part). As a consequence,

$$\mathbf{a} \rightarrow \mathbf{a}'' = \mathbf{v}\mathbf{u}\mathbf{a}\mathbf{u}\mathbf{v}$$

is the **rotation** of \mathbf{a} in the plane $\mathbf{u} \wedge \mathbf{v}$ through twice the angle between \mathbf{u} and \mathbf{v} .

Vibrational part of metric tensor

- The vibrational metric tensor elements is obtained as

$$g^{(s_i s_j)} = \sum_{\alpha=1}^N \frac{(\nabla_{\mathbf{x}_\alpha} s_i) \cdot (\nabla_{\mathbf{x}_\alpha} s_j)}{m_\alpha} \quad (4)$$

- By utilizing geometric algebra, it is easy to evaluate the gradients $\nabla_{\mathbf{x}_\alpha} s_i$.

Shape coordinate gradients

Example

Gradient $\nabla_{\mathbf{x}_\beta} r_{\alpha\beta}$ of **bond length** $r_{\alpha\beta} = |\mathbf{r}_{\alpha\beta}| = (\mathbf{r}_{\alpha\beta} \cdot \mathbf{r}_{\alpha\beta})^{1/2}$ (now $\mathbf{r}_{\alpha\beta} = \mathbf{x}_\beta - \mathbf{x}_\alpha$). Start from

$$r_{\alpha\beta}^2 = \mathbf{r}_{\alpha\beta} \cdot \mathbf{r}_{\alpha\beta}$$

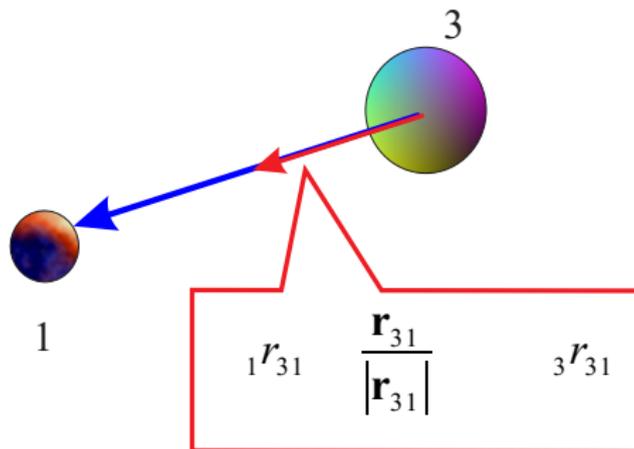
Take the vectorial derivative of both sides as

$$2r_{\alpha\beta} \nabla_{\mathbf{x}_\beta} r_{\alpha\beta} = 2\mathbf{r}_{\alpha\beta}$$

to get

$$\nabla_{\mathbf{x}_\beta} r_{\alpha\beta} = \frac{\mathbf{r}_{\alpha\beta}}{r_{\alpha\beta}} = -\nabla_{\mathbf{x}_\alpha} r_{\alpha\beta}$$

Shape coordinate gradients (cntd.)



Rotational and Coriolis parts of metric tensor

- Rotational and Coriolis parts of the metric tensor elements are given by

$$g^{(l'_i l'_j)} = \sum_{\alpha=1}^N \frac{\mathbf{e}_{\alpha}^{(l'_i)} \cdot \mathbf{e}_{\alpha}^{(l'_j)}}{m_{\alpha}} \quad (5)$$

$$g^{(l'_i s_j)} = \sum_{\alpha=1}^N \frac{\mathbf{e}_{\alpha}^{(l'_i)} \cdot (\nabla_{\mathbf{x}_{\alpha}} s_j)}{m_{\alpha}} \quad (6)$$

- $\mathbf{e}_{\alpha}^{(l'_k)}$ is the **measuring vector** associated to nucleus α and k th body-frame component of the angular momentum operator.

Rotational measuring vectors

- Generally

$$(\nabla_{\mathbf{x}_\alpha} \theta) p_\theta + (\nabla_{\mathbf{x}_\alpha} \chi) p_\chi + (\nabla_{\mathbf{x}_\alpha} \phi) p_\phi = \sum_k^3 \mathbf{e}_\alpha^{(l'_k)} l'_k \quad (7)$$

- However,

$$\mathbf{e}_\alpha^{(l'_i)} \neq \nabla_{\mathbf{x}_\alpha} \theta, \nabla_{\mathbf{x}_\alpha} \chi, \nabla_{\mathbf{x}_\alpha} \phi \quad (8)$$

i.e. components of angular momentum **are not conjugated to any rotational coordinates**.

- By a short calculation (omitted here), it turns out that

$$\mathbf{e}_\alpha^{(l'_i)} = \nabla_{\mathbf{a}} [(\mathbf{a} \cdot \nabla_{\mathbf{x}_\alpha} \mathbf{u}'_j) \cdot \mathbf{u}'_k] \quad (9)$$

where i, j, k are in cyclic order.

Rotational measuring vectors (cntnd.)

Example

Let us choose

$$\mathbf{u}'_3 = \hat{\mathbf{r}}_{12}$$

Then

$$\mathbf{a} \cdot \nabla_{\mathbf{x}_2} \mathbf{u}'_3 = \mathbf{a} \cdot \nabla_{\mathbf{x}_2} \hat{\mathbf{r}}_{12} = \frac{\mathbf{a} - \mathbf{a} \cdot \hat{\mathbf{r}}_{12} \hat{\mathbf{r}}_{12}}{|\mathbf{r}_{12}|}$$

Hence,

$$\mathbf{e}_2^{(l'2)} = \nabla_{\mathbf{a}} [(\mathbf{a} \cdot \nabla_{\mathbf{x}_2} \mathbf{u}'_3) \cdot \mathbf{u}'_1] = \nabla_{\mathbf{a}} \frac{\mathbf{a} \cdot \mathbf{u}'_1}{|\mathbf{r}_{12}|} = \frac{\mathbf{u}'_1}{|\mathbf{r}_{12}|}$$

Jacobian

- Conventionally the **Jacobian** J is evaluated as $J = (\det g)^{-1/2}$.
In practice, **way** too complicated!
- Better, $J = J_1 J_2 J_3 \dots$, where

$$J_1 = \frac{\sin \theta}{\left| \mathbf{e}_1^{(l'_1)} \wedge \mathbf{e}_1^{(l'_2)} \wedge (\nabla_{\mathbf{x}_1} s_1) \right|}$$

$$J_2 = \frac{1}{\left| \mathbf{e}_2^{(l'_3)} \wedge (\nabla_{\mathbf{x}_2} s_2) \wedge (\nabla_{\mathbf{x}_2} s_3) \right|}$$

$$J_\alpha = \frac{1}{\left| (\nabla_{\mathbf{x}_\alpha} s_i) \wedge (\nabla_{\mathbf{x}_\alpha} s_j) \wedge (\nabla_{\mathbf{x}_\alpha} s_k) \right|} \text{ for } \alpha = 3, 4, \dots, N-1$$

Jacobian (cntnd.)

Example

Consider **polyspherical** coordinates. The jacobian reads as

$$J = \sin \theta \prod_{i=1}^{N-1} r_i^2 \prod_{j=1}^{N-2} \sin \theta_j \quad (10)$$

where the bond lengths have been denoted as r_i , the bond-angles as θ_j , and the torsional (dihedral) angles as ϕ_k .

Recent advances

- GA approach produces KEOs also when the implicitly defined **Eckart body-frames** are used.
 - Eckart frames occupy very important role in the development of vibration-rotation theory.
- **Constraints in shape** do not cause problems.
 - A correction term is added to each active measuring vectors.
 - Riemannian quantization.

Conclusions

- Geometric algebra is an ideal tool for obtaining vibration-rotation KEOs
 - Applies for arbitrary coordinates and body-frames.
 - Based on measuring vectors.

Further Reading



J. Pesonen.

Vibrational coordinates and their gradients: A geometric algebra approach.

Journal of Chemical Physics, **112**(7):3121, 2000.



J. Pesonen.

Vibration–rotation kinetic energy operators: A geometric algebra approach.

Journal of Chemical Physics, **114**(24):10598, 2001.



J. Pesonen.

Volume-elements of integration: A geometric algebra approach.

Journal of Chemical Physics, **116**(24):1825, 2002.